

When Non-Expanding Horizons met Near Horizon Geometries

Jerzy Lewandowski, Adam Szereszewski, Piotr Waluk
Uniwersytet Warszawski

A long time ago in a galaxy far, far away....

A peculiar solution to Einstein's Equations was constructed

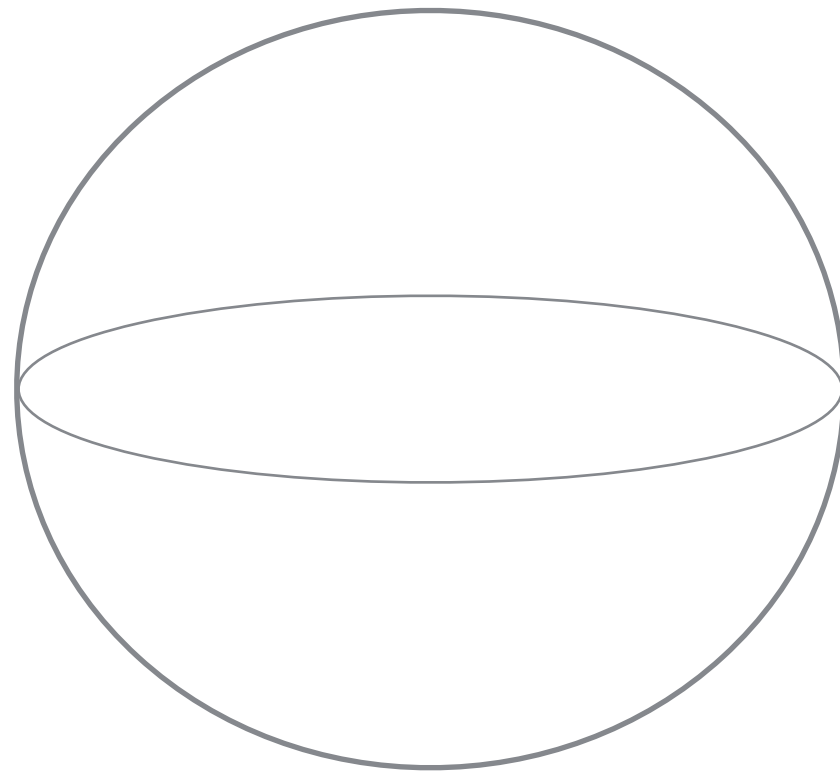
Pawłowski, Lewandowski, Jezierski - 2003

Spacetime foliated by the Killing horizons ...

Ingredients

$$S^2, \quad g_{AB}, \quad \omega_A$$

a 2-sphere,
a metric tensor, a 1-form



Such that :

$${}^2\nabla_A \omega_B + {}^2\nabla_B \omega_A + 2\omega_A \omega_B - {}^2R_{AB} = 0$$

*Ashtekar, Beetle, JL 2001; JL, Pawłowski 2002;
Chruściel, Reall, Tod 2005; Jezierski 2008, Jezierski, Kamiński 2012*

$$S^2, \quad g_{AB}, \quad \omega_A$$

$${}^2\nabla_A \omega_B + {}^2\nabla_B \omega_A + 2\omega_A \omega_B - {}^2R_{AB} = 0$$

The solution

On: $S^2 \times \mathbb{R} \times \mathbb{R}$
 parametrized by: (x^A, u, v)

$$g_{\mu\nu} dx^\mu dx^\nu :=$$

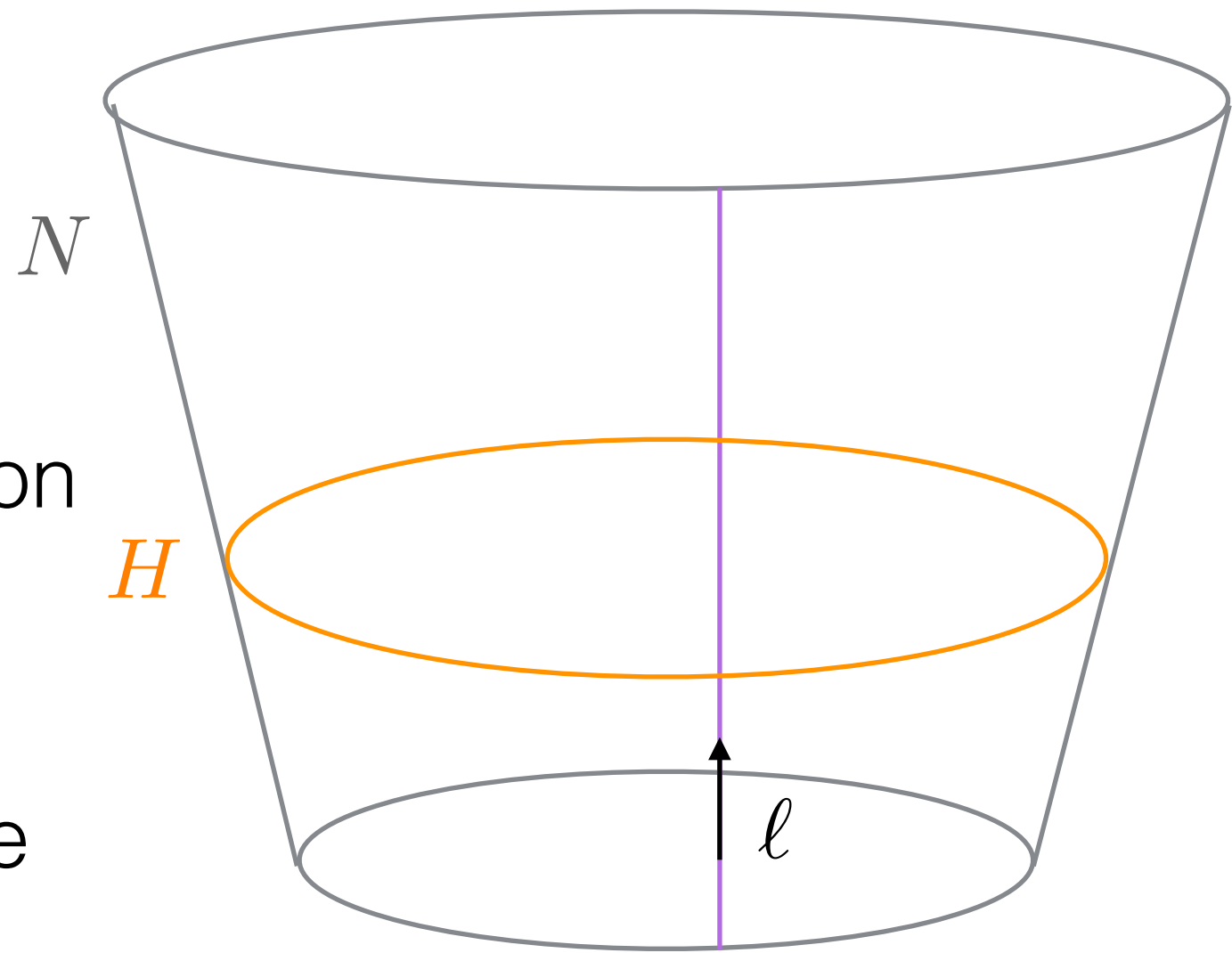
$$g_{AB} dx^A dx^B - 2du \left(dv - 2v\omega - \frac{1}{2}v^2(\operatorname{div}\omega + 2\omega^2)du \right)$$

$$R_{\mu\nu} = 0$$

The present day

Non-expanding horizons: definition

- a codimension 1
- null surface
- non-expanding
- compact space-like section



Assumptions about spacetime

$$M, g_{\mu\nu} \quad - + \dots +$$

$$\dim M = n \geq 3$$

$$G_{\mu\nu} = T_{\mu\nu}$$

$$\ell^\mu \ell_\mu = 0$$

$$T_{\ell\ell} \geq 0$$

$$T_{\ell\alpha} T_\ell{}^\alpha \leq 0$$

Higher dimensional Raychaudhuri equation

JL, Pawłowski 2004

Every null, geodesic vector field tangent to a null surface satisfies:

$$\ell^\mu \nabla_\mu \theta = -\frac{1}{n-2} \theta^2 - \sigma_{AB} \sigma^{AB} - T_{\ell\ell}$$

Therefore:

$$\theta = 0 \quad \Rightarrow \quad \mathcal{L}_\ell g_{ab} = 0$$

Non-expanding horizons: geometry

$g_{\mu\nu}, \nabla_\mu$ - spacetime geometry

$\mu, \nu - M$

$a, b - N$

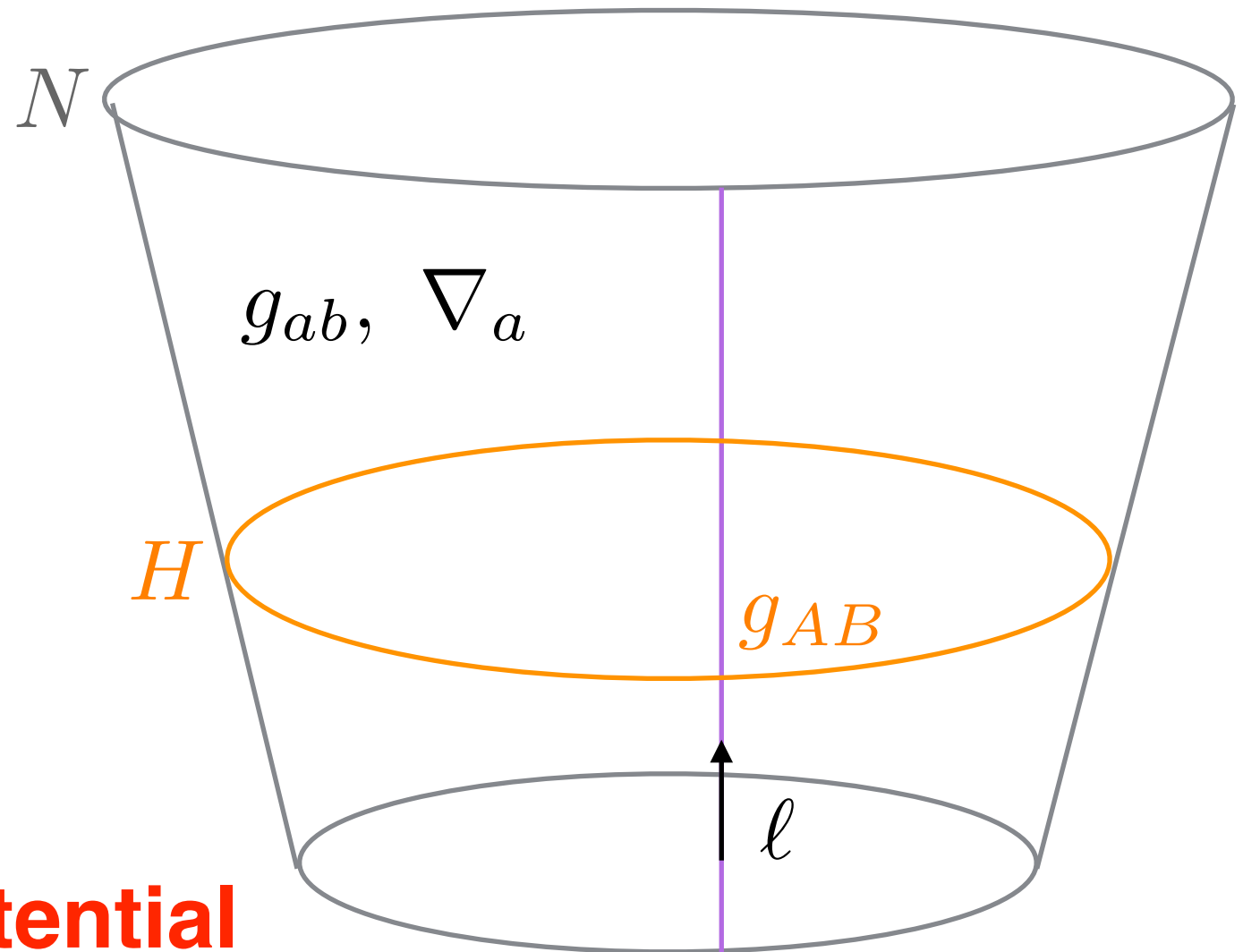
$A, B - H$

g_{ab}, ∇_a - **horizon geometry**

$\nabla_a \ell^b = \omega_a^\ell \ell^b$ - **rotation potential**

$\kappa^\ell := \ell^a \omega_a^\ell$ - **surface gravity**

$\nabla_a \kappa^\ell - \mathcal{L}_\ell \omega_a^\ell = 0$ - **the zeroth law**



Non-expanding horizons: symmetric

Symmetric NEH:

$$\begin{aligned}\mathcal{L}_X g_{ab} &= 0 \\ [\mathcal{L}_X, \nabla_a] &= 0\end{aligned}$$

Isolated Horizon:

$$[\mathcal{L}_\ell, \nabla_a] = 0$$

$$\begin{aligned}\mathcal{L}_\ell \omega_a^\ell &= 0 \\ \kappa^\ell &= \text{const}\end{aligned}$$

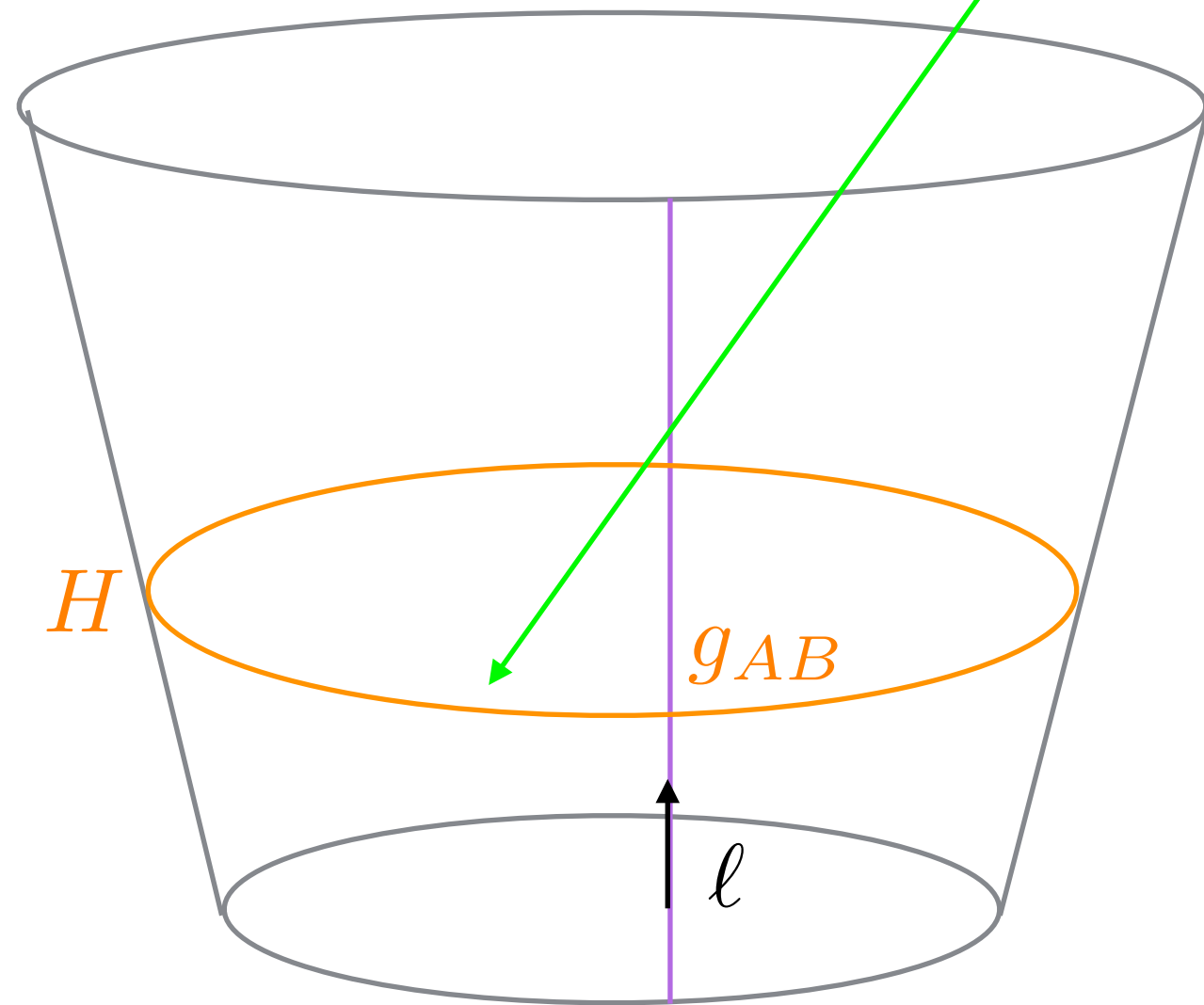
Extremal Isolated Horizons

$$\kappa^\ell = 0$$

$$(1) \quad (n-2)\nabla_A \omega_B^\ell + (n-2)\nabla_B \omega_A^\ell + 2\omega_A^\ell \omega_B^\ell - (n-2)R_{AB} + \overset{\substack{\text{red arrow} \\ T_{\mu\nu}}}{R_{AB}} = 0$$

The constraint satisfied
on every $n-2$ -slice

**Conversely, if a
non-expanding horizon,
admits a slice and a
null vector such that
the condition (1) is
satisfied, then it is an
extremal isolated
horizon provided:**



$$\mathcal{L}_\ell R_{ab} = 0$$

References to the extremal isolated horizons:

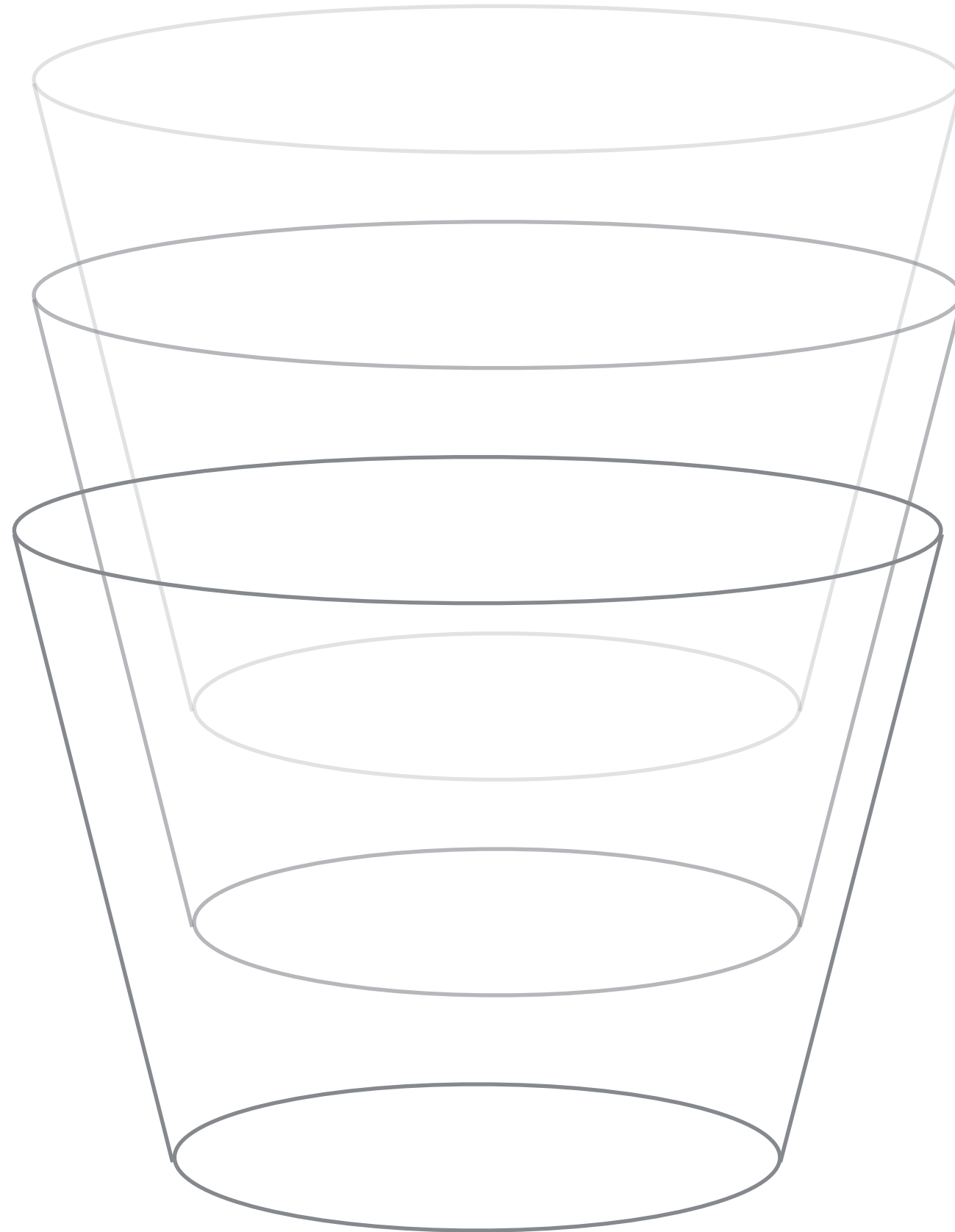
Ashtekar, JL, Beetle 2001;

JL, Pawłowski 2002;

JL, Pawłowski 2004;

JL, Szereszewski, Waluk 2016

Spacetimes foliated by non-expanding horizons



Spacetimes foliated by non-expanding horizons

Let

$$u = \text{const}$$

define the foliation.

$$\ell' = -g^{\mu\nu} \nabla_\nu u$$

On each space-like section of every non-expanding horizon the following condition is satisfied

$$-(n-2) \nabla_A \omega_B^{\ell'} - (n-2) \nabla_B \omega_A^{\ell'} + 2 \omega_A^{\ell'} \omega_B^{\ell'} - (n-2) R_{AB} + R_{AB} = 0$$

Compare the very condition satisfied by foliating horizons

$$-(n-2)\nabla_A\omega_B^{\ell'} - (n-2)\nabla_B\omega_A^{\ell'} + 2\omega_A^{\ell'}\omega_B^{\ell'} - (n-2)R_{AB} + R_{AB} = 0$$

with the extremal horizon condition

$$(n-2)\nabla_A\omega_B^{\ell} + (n-2)\nabla_B\omega_A^{\ell} + 2\omega_A^{\ell}\omega_B^{\ell} - (n-2)R_{AB} + R_{AB} = 0$$

The transformation

$$\omega^{\ell} = -\omega^{\ell'}$$

would do. But where from?

Intersecting non-expanding horizons

Suppose two horizons intersect:

$$N \cap N' = H$$

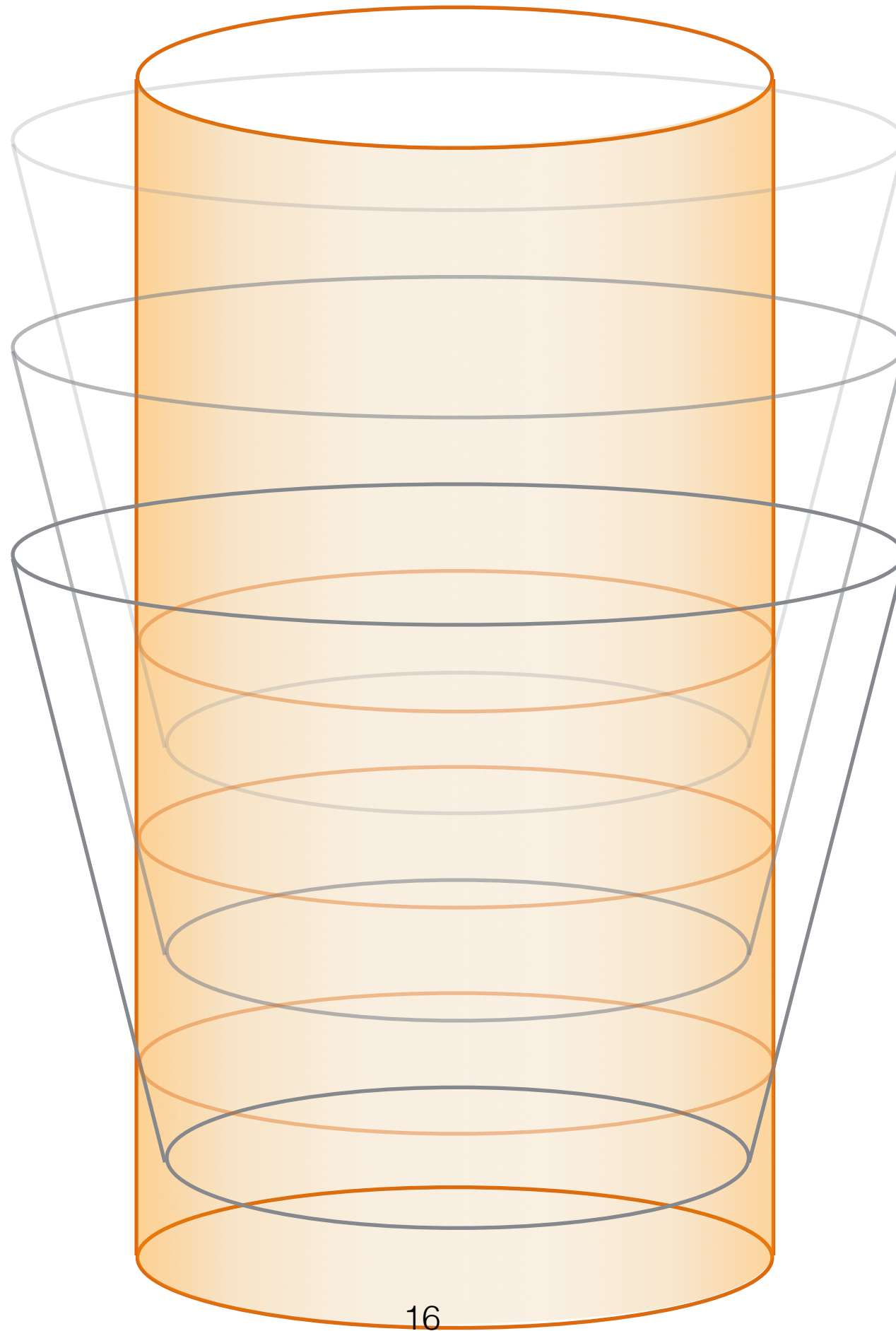
and

$$\ell^\mu \ell'_\mu = -1$$

Then

$$\omega_A^\ell = -\ell'_\mu \nabla_A \ell^\mu = \ell_\mu \nabla_A \ell'^\mu = -\omega_A^{\ell'}$$

Non-expanding horizon foliation and a transversal horizon



Results valid for vacuum

1. Suppose spacetime is foliated by non-expanding horizons and there exists a transversal non-expanding horizon. Then, the transversal horizon is extremal. Still true with non-zero cosmological constant.

2. Given two intersecting non-expanding horizons, the Cauchy problem is well posed given

$$g_{AB}, \omega_{AB}$$

on the intersection, hence this data determines a solution

3. A general solution to 1 above is:

Given an $n-1$ dimensional manifold, a metric and a 1-form thereon:

$$H, \quad g_{AB}, \quad \omega_A$$

Such that:

$$^{(n-2)}\nabla_A \omega_B + ^{(n-2)}\nabla_B \omega_A + 2\omega_A \omega_B - ^{(n-2)}R_{AB} = 0$$

On:

$$H \times \mathbb{R} \times \mathbb{R} \\ (x^A, u, v)$$

Define:

$$g_{\mu\nu} dx^\mu dx^\nu := \\ g_{AB} dx^A dx^B - 2du \left(dv - 2v\omega - \frac{1}{2}v^2(\operatorname{div}\omega + 2\omega^2)du \right)$$

$$R_{\mu\nu} = 0$$

The Spacetime Killing vectors

$$g_{AB}dx^A dx^B = 2du \left(dv - 2v\omega - \frac{1}{2}v^2(\text{div}\omega + 2\omega^2)du \right)$$

Killing vectors:

$$K = \partial_u \qquad \xi = v\partial_v - u\partial_u$$

Killing horizons of $u_0 K + \xi$

$$u = u_0 \quad \text{or} \quad v = 0$$

Extremal Killing horizon of K

$$v = 0$$

Near horizon geometries

Geometry in a neighborhood of an extremal Killing horizon

$$v = 0$$

$$g_{AB}(x, v)dx^A dx^B - 2du (dv + vW_A(x, v)dx^A + v^2 H(x, v)du)$$

Horowitz transformation

$$v \mapsto \epsilon v, \quad u \mapsto \frac{u}{\epsilon}$$

Limit at $\epsilon \mapsto 0$ **a near horizon geometry:**

$$g_{AB}(x, 0)dx^A dx^B - 2du(dv + vW_A(x, 0)dx^A + v^2 H(x, 0)du)$$

Horowitz 1999, Reall 2002; Kunduri, Lucietti 2013

Summary

1. We investigated n dimensional spacetimes foliated by non-expanding horizons. We found out that:
 - a) each foliating horizon satisfies a constraint dual to the extremal isolated horizon constraint
 - b) the duality can be geometrically interpreted as a condition allowing (but not implying) the existence of a transversal extremal horizon.
 - c) if we assume that there actually exists a transversal extremal horizon, than we can solve vacuum Einstein's equation completely.
 - d) The general solution to our problem (foliation by NEH and a transversal NEH) turned out to be a general **near horizon geometry**.

Summary

- e) In the $n=4$ case we derived all them already in the earlier paper by *Pawłowski, JL, Jezierski of 2003*
- f) Remarkably in the two major Living Reviews: *Ashtekar, Krishnan 2004 Isolated and Dynamical Horizons* our paper is mentioned as “another exotic case”
Kunduri, Lucietti Living Reviews 2013 Near Horizon Geometries the paper is unnoticed

2. Open questions:

- a) Are there other solutions?
- b) Does our “master equation” on 2-sphere have other solutions than those defined by the extremal Kerr? If you find a solution you have a new exact solution to Einsteins Equations.

Thank You